

MAT 576

Fall 2008

Homework set 4 Answers

1. In class we looked at the wave equation $u_{xx} - u_{yy} = 0$ on a bounded domain $\Omega \subset \mathbb{R}^2$ that is partitioned into two (open) subdomains Ω_1 to the left of a smooth curve $\mathcal{C} : x = s(t)$ and Ω_2 to the right. Letting $[[w]] := w|_{\Omega_2} - w|_{\Omega_1}$ we saw that a function u that satisfies $u_{xx} - u_{yy} = 0$ in $\Omega_1 \cup \Omega_2$ is a weak solution only if for all $\phi \in C_0^\infty(\Omega)$ we have

$$\int_{\mathcal{C}} ([[u_x]] N_1 \phi - [[u_y]] N_2 \phi) dl + \int_{\mathcal{C}} ([[u]] N_1 \phi_x - [[u]] N_2 \phi_y) dl = 0,$$

where $\langle N_1, N_2 \rangle$ is the (continuously varying) unit normal to \mathcal{C} . Show that the function $H(x - y)$ satisfies this condition.

2. This is an exercise in using Green's theorem (also known as Gauss's theorem or Stokes's theorem, or Ostrogradskii's theorem or by various combinations of these names):

$$\int_{\Omega} u_{x_i} v dx = - \int_{\Omega} uv_{x_i} dx + \int_{\partial\Omega} uv N_i dS,$$

where $N := \langle N_1, \dots, N_n \rangle$ is the unit outward normal to $\partial\Omega$ and Ω a bounded domain in \mathbb{R}^n , $n > 1$, with piecewise smooth boundary. Let $u \in C^2(\Omega) \cap C^1(\bar{\Omega})$ and suppose that u is a solution to the nonlinear elliptic problem

$$\sum_{i=1}^n \frac{\partial^2 u}{\partial x_i^2} = f(u) \text{ in } \Omega, \quad u = 0 \text{ on } \partial\Omega.$$

Suppose that f is continuous and define $F(s) = \int_0^s f(\sigma) d\sigma$. We employ the summation convention.

- (a) Show that

$$\int_{\Omega} x_j \frac{\partial^2 u}{\partial x_j \partial x_i} \frac{\partial u}{\partial x_i} dx = -\frac{n}{2} \int_{\Omega} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} dx + \frac{1}{2} \int_{\partial\Omega} (N \cdot x) \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} dS.$$

- (b) Show that

$$\int_{\Omega} f(u) x_i \frac{\partial u}{\partial x_i} dx = \int_{\partial\Omega} F(u) x_i N_i dS - n \int_{\Omega} F(u) dx.$$

- (c) Show that

$$\frac{n-2}{2} \int_{\Omega} \|\nabla u\|^2 dx + \frac{1}{2} \int_{\partial\Omega} \left(\frac{\partial u}{\partial N} \right)^2 (x \cdot N) dS + n \int_{\Omega} F(u) dx = 0.$$

Here $\partial u / \partial N = \nabla u \cdot N$ is the outward normal derivative. Hint: use the fact that $u = 0$ on the boundary - at least twice!

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Let Ω be a bounded domain in \mathbb{R}^n , $n > 1$, with piecewise smooth boundary. Let $N := (N_1, \dots, N_n)^T$ be the unit outward normal to $\partial\Omega$. It should be noted that N exists almost everywhere on the boundary by smoothness, and so boundary integrals involving N make sense. Let $u \in C^2(\Omega) \cap C^1(\bar{\Omega})$ be a solution to the nonlinear elliptic problem

$$\sum_{i=1}^n \frac{\partial^2 u}{\partial x_i^2} = f(u) \text{ in } \Omega, \quad u = 0 \text{ on } \partial\Omega$$

where f is continuous. Define $F(s) := \int_0^s f(\sigma) d\sigma$. In what follows the summation convention is observed on latin indices, but greek indices are not summed. Both run from 1 to n .

(a) For fixed α ,

$$\int_{\Omega} x_{\alpha} \frac{\partial^2 u}{\partial x_{\alpha} \partial x_i} \frac{\partial u}{\partial x_i} dx = \int_{\Omega} x_{\alpha} \left(\frac{1}{2} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} \right)_{x_{\alpha}} dx \quad (1)$$

$$= \frac{1}{2} \int_{\partial\Omega} x_{\alpha} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} N_{\alpha} dS - \frac{1}{2} \int_{\Omega} (x_{\alpha})_{x_{\alpha}} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} dx \quad (2)$$

$$= \frac{1}{2} \int_{\partial\Omega} x_{\alpha} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} N_{\alpha} dS - \frac{1}{2} \int_{\Omega} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} dx \quad (3)$$

where equality (2) is by Green's Theorem. Summing over α and noting that the second integrand does not depend on α yields the result

$$\int_{\Omega} x_j \frac{\partial^2 u}{\partial x_j \partial x_i} \frac{\partial u}{\partial x_i} dx = -\frac{n}{2} \int_{\Omega} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} dx + \frac{1}{2} \int_{\partial\Omega} (x \cdot N) \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} dS$$

(b) Let

$$F(s) = \int_0^s f(\sigma) d\sigma$$

So,

$$F'(s) = f(s)$$

For fixed α ,

$$\int_{\Omega} F(u) dx = \int_{\Omega} (x_{\alpha})_{x_{\alpha}} F(u) dx \quad (4)$$

$$= \int_{\partial\Omega} x_{\alpha} F(u) N_{\alpha} dS - \int_{\Omega} x_{\alpha} (F(u))_{x_{\alpha}} dx \quad (5)$$

$$= \int_{\partial\Omega} x_{\alpha} F(u) N_{\alpha} dS - \int_{\Omega} x_{\alpha} F'(u) \frac{\partial u}{\partial x_{\alpha}} dx \quad (6)$$

$$= \int_{\partial\Omega} x_{\alpha} F(u) N_{\alpha} dS - \int_{\Omega} x_{\alpha} \frac{\partial u}{\partial x_{\alpha}} f(u) dx \quad (7)$$

where equality (5) is by Green's Theorem, (6) is by the chain rule, and (7) is by our earlier computation. Summing over α , noting that the first integrand does not depend on α and rearranging yields the result

$$\int_{\Omega} f(u) x_i \frac{\partial u}{\partial x_i} dx = \int_{\partial\Omega} F(u) x_i N_i dS - n \int_{\Omega} F(u) dx$$

(c) Since $u = 0$ on $\partial\Omega$, $F(u) = F(0) = 0$ on $\partial\Omega$. Also, because Ω is open in \mathbb{R}^n , $\partial\Omega$ is an $(n-1)$ -dimensional piecewise smooth hypersurface in \mathbb{R}^n . Thus the tangent space at a non-cusp point of $\partial\Omega$ is dimension $n-1$. Hence at a non-cusp point $x \in \partial\Omega$ there exists a basis for \mathbb{R}^n which decomposes \mathbb{R}^n as the orthogonal direct sum $\mathbb{R}^n = T_x(\partial\Omega) \oplus N$, where $T_x(\partial\Omega)$ is the tangent space to $\partial\Omega$ at x . Now, since $u = 0$ on $\partial\Omega$, the derivative of u in any direction tangent to $\partial\Omega$ is 0. It thus follows from the direct-sum decomposition above that $\nabla u = aN$ for some scalar a . In fact, $\nabla u = \frac{\partial u}{\partial N} N$, and so for any i ,

$$\frac{\partial u}{\partial x_i} = \frac{\partial u}{\partial N} N_i$$

For fixed α ,

$$\int_{\Omega} x_i \frac{\partial u}{\partial x_i} \frac{\partial^2 u}{\partial x_{\alpha} \partial x_{\alpha}} dx = \int_{\Omega} x_i \frac{\partial u}{\partial x_i} \left(\frac{\partial u}{\partial x_{\alpha}} \right)_{x_{\alpha}} dx \quad (8)$$

$$= \int_{\partial\Omega} x_i \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_{\alpha}} N_{\alpha} dS - \int_{\Omega} \frac{\partial u}{\partial x_{\alpha}} \left(x_i \frac{\partial u}{\partial x_i} \right)_{x_{\alpha}} dx \quad (9)$$

$$= \int_{\partial\Omega} x_i \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_{\alpha}} N_{\alpha} dS - \int_{\Omega} \left(x_i \frac{\partial u}{\partial x_{\alpha}} \frac{\partial^2 u}{\partial x_i \partial x_{\alpha}} + \frac{\partial u}{\partial x_{\alpha}} \frac{\partial u}{\partial x_{\alpha}} \right) dx \quad (10)$$

where equality (10) is by Green's Theorem and (11) is by the chain rule. Summing over α yields the result

$$\int_{\Omega} x_i \frac{\partial u}{\partial x_i} \frac{\partial^2 u}{\partial x_j \partial x_j} dx = \int_{\partial\Omega} x_i \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} N_j dS - \int_{\Omega} \left(x_i \frac{\partial u}{\partial x_j} \frac{\partial^2 u}{\partial x_i \partial x_j} + \frac{\partial u}{\partial x_j} \frac{\partial u}{\partial x_j} \right) dx$$

Now,

$$n \int_{\Omega} F(u) dx = \int_{\partial\Omega} F(u) x_i N_i dS - \int_{\Omega} f(u) x_i \frac{\partial u}{\partial x_i} dx \quad (11)$$

$$= - \int_{\Omega} x_i \frac{\partial u}{\partial x_i} \frac{\partial^2 u}{\partial x_j \partial x_j} dx \quad (12)$$

where (13) is by (b) and (14) is by $u = 0$ on $\partial\Omega$ and the PDE. From our above calculation we obtain

$$= - \int_{\partial\Omega} x_i \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} N_j dS + \int_{\Omega} x_i \frac{\partial u}{\partial x_j} \frac{\partial^2 u}{\partial x_i \partial x_j} dx + \int_{\Omega} \frac{\partial u}{\partial x_j} \frac{\partial u}{\partial x_j} dx \quad (13)$$

$$= - \int_{\partial\Omega} x_i \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_j} N_j dS + \int_{\Omega} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} dx - \frac{n}{2} \int_{\Omega} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} dx + \frac{1}{2} \int_{\partial\Omega} (x \cdot N) \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} dS \quad (14)$$

$$= \frac{2-n}{2} \int_{\Omega} \frac{\partial u}{\partial x_i} \frac{\partial u}{\partial x_i} dx - \int_{\partial\Omega} \left(\frac{\partial u}{\partial N} \right) N_i \left(\frac{\partial u}{\partial N} \right) N_j N_j dS + \frac{1}{2} \int_{\partial\Omega} x_i N_i \left(\frac{\partial u}{\partial N} \right)^2 N_j N_j dS \quad (15)$$

$$= \frac{2-n}{2} \int_{\Omega} \|\nabla u\|^2 dx - \frac{1}{2} \int_{\partial\Omega} \left(\frac{\partial u}{\partial N} \right)^2 (x \cdot N) dS \quad (16)$$

where (15) is by (12), (16) is by (a), and (17) is by (8). Rearranging yields the result

$$\frac{n-2}{2} \int_{\Omega} \|\nabla u\|^2 dx + \frac{1}{2} \int_{\partial\Omega} \left(\frac{\partial u}{\partial N} \right)^2 (x \cdot N) dS + n \int_{\Omega} F(u) dx = 0$$

3. Consider the conservation law with 2 components (1 space variable) $\mathbf{u}_t + \mathbf{f}(\mathbf{u})_x = \mathbf{0}$, or equivalently $\mathbf{u}_t + \mathbf{A}(\mathbf{u})\mathbf{u}_x = \mathbf{0}$. Let $\lambda_1(\mathbf{u})$ and $\lambda_2(\mathbf{u})$ be the eigenvalues of $\mathbf{A}(\mathbf{u})$ with respective eigenvectors \mathbf{r}_1 and \mathbf{r}_2 and let $w_1(\mathbf{u})$ and $w_2(\mathbf{u})$ be corresponding Riemann invariants. Let $\mathbf{v}(x, t) := (v^1(x, t), v^2(x, t)) = \mathbf{w}(\mathbf{u}(x, t))$. Show that

$$v_t^1 + \lambda_2(\mathbf{u})v_x^1 = 0$$

$$v_t^2 + \lambda_1(\mathbf{u})v_x^2 = 0$$

4. Consider the conservation law with 2 components (1 space variable) $\mathbf{u}_t + \mathbf{f}(\mathbf{u})_x = \mathbf{0}$, or equivalently $\mathbf{u}_t + \mathbf{A}(\mathbf{u})\mathbf{u}_x = \mathbf{0}$. Let $\lambda_1(\mathbf{u})$ and $\lambda_2(\mathbf{u})$ be the eigenvalues of $\mathbf{A}(\mathbf{u})$ with respective eigenvectors \mathbf{r}_1 and \mathbf{r}_2 and let $w_1(\mathbf{u})$ and $w_2(\mathbf{u})$ be corresponding Riemann invariants. Let $\mathbf{v}(x, t) := (v^1(x, t), v^2(x, t)) = \mathbf{w}(\mathbf{u}(x, t))$. Show that

$$v_t^1 + \lambda_2(\mathbf{u})v_x^1 = 0$$

$$v_t^2 + \lambda_1(\mathbf{u})v_x^2 = 0$$

5. One-dimensional isentropic flow of a gas is modeled by the system

$$\rho_t + v\rho_x + \rho v_x = 0$$

$$v_t + vv_x + c^2\rho_x/\rho = 0,$$

where c , which is positive, is the speed of sound and is a function of the density ρ , $c^2 = K\rho^{\gamma-1}$, $\gamma > 1$ is the ratio of the heat capacities (v is the velocity of the flow).

(a) Write this system as a 2-dimensional conservation law $\mathbf{u}_t + \mathbf{f}(\mathbf{u})_x = \mathbf{0}$ for $\mathbf{u} = (\rho, v)^T$.

(b) Write the system as $\mathbf{u}_t + \mathbf{A}\mathbf{u}_x = \mathbf{0}$ and find the eigenvalues and eigenvectors of $\mathbf{A}(\mathbf{u})$.

- (c) Introduce the function $L(\rho) = \int_0^\rho (c(\sigma)/\sigma) d\sigma$ and find Riemann invariants $q(\mathbf{u})$ corresponding to the eigenvalue λ_1 and $s(\mathbf{u})$ corresponding to the eigenvalue λ_2 . Pick the labeling so that $\lambda_1 > \lambda_2$.
- (d) Use problem 3 to find the system of equations for q and s .
- (e) Show that the system can be “inverted” to obtain:

$$\begin{aligned}x_q &= (v - c)t_q \\x_s &= (v + c)t_s\end{aligned}$$

Hint:

$$\begin{pmatrix} \frac{\partial x}{\partial q} & \frac{\partial x}{\partial s} \\ \frac{\partial t}{\partial q} & \frac{\partial t}{\partial s} \end{pmatrix} = \begin{pmatrix} \frac{\partial q}{\partial x} & \frac{\partial q}{\partial t} \\ \frac{\partial s}{\partial x} & \frac{\partial s}{\partial t} \end{pmatrix}^{-1} = \frac{1}{D} \begin{pmatrix} \frac{\partial s}{\partial t} & -\frac{\partial q}{\partial t} \\ -\frac{\partial s}{\partial x} & \frac{\partial q}{\partial x} \end{pmatrix}, \quad D = \frac{\partial(q, s)}{\partial(x, t)}.$$

- (f) Show that $v + c$ and $v - c$ can be written as functions of q and s and that the last system is therefore a fully linear system.

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$$\begin{pmatrix} \rho \\ v \end{pmatrix}_t + \begin{pmatrix} \rho v \\ \frac{1}{2}v^2 + \frac{K\rho^{\gamma-1}}{\gamma-1} \end{pmatrix}_x = 0$$

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$$\begin{pmatrix} \rho \\ v \end{pmatrix}_t + \begin{pmatrix} v & \rho \\ c^2/\rho & v \end{pmatrix} \begin{pmatrix} \rho \\ v \end{pmatrix}_x = 0$$

- The characteristic equation for the above matrix is $\lambda^2 - 2v\lambda + v^2 - c^2 = 0$. This yields $\lambda_1 = v + c$ and $\lambda_2 = v - c$. The corresponding eigenvectors are $(\rho, c)^T$ and $(\rho, -c)^T$. We want $(q_\rho, q_v)^T \perp (\rho, c)$. This is easier to figure out if we separate the v and ρ in the eigenvector; that is to say, use the eigenvector $(1, c/\rho)$ instead. We then see that $q := v - L(\rho)$ works. Similarly, for λ_2 we get the Riemann invariant $s := v + L(\rho)$.
- From problem 3 we immediately get

$$\begin{aligned}q_t + (v - c)q_x &= 0, \\s_t + (v + c)s_x &= 0.\end{aligned}$$

- Since

$$\begin{pmatrix} D \frac{\partial x}{\partial q} & D \frac{\partial x}{\partial s} \\ D \frac{\partial t}{\partial q} & D \frac{\partial t}{\partial s} \end{pmatrix} = \begin{pmatrix} \frac{\partial s}{\partial t} & -\frac{\partial q}{\partial t} \\ -\frac{\partial s}{\partial x} & \frac{\partial q}{\partial x} \end{pmatrix},$$

we can rewrite the system as

$$\begin{aligned}x_q &= (v + c)t_q, \\x_s &= (v - c)t_s.\end{aligned}$$

- Using the fact that $q = v - 2c/(\gamma - 1)$ and $s = v + 2c/(\gamma - 1)$ we get $v = (q + s)/2$ and $c = (\gamma - 1)(s - q)/4$. A little algebra then yields

$$\begin{aligned}v + c &= \frac{(3 - \gamma)q + (1 + \gamma)s}{4}, \\v - c &= \frac{(3 - \gamma)s + (1 + \gamma)q}{4}.\end{aligned}$$